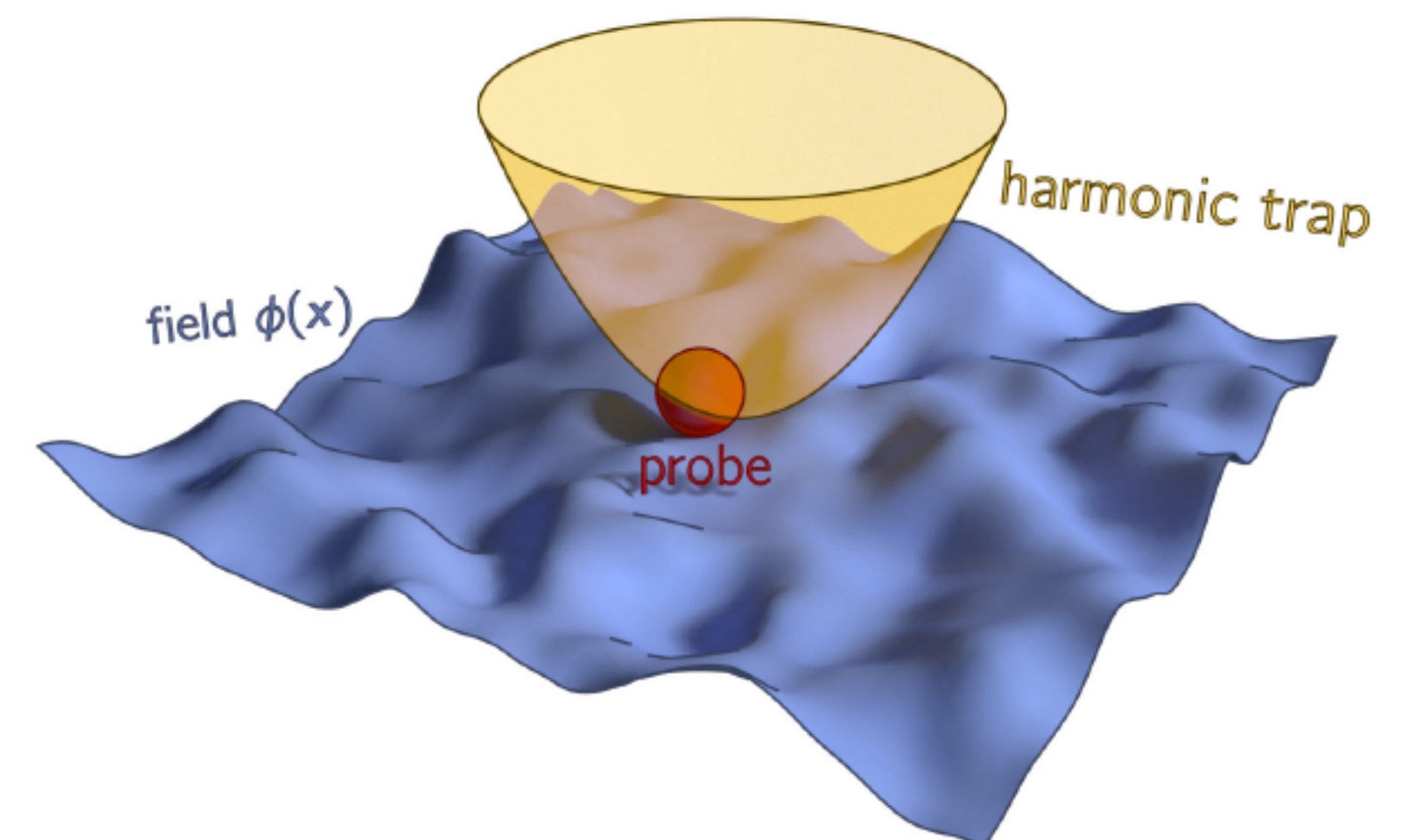


Dynamics of tracer particles in a fluctuating correlated medium

Davide Venturelli

StatPhys28, Tokyo, August 9th 2023



Particle in a complex medium

From Brownian motion to correlated media

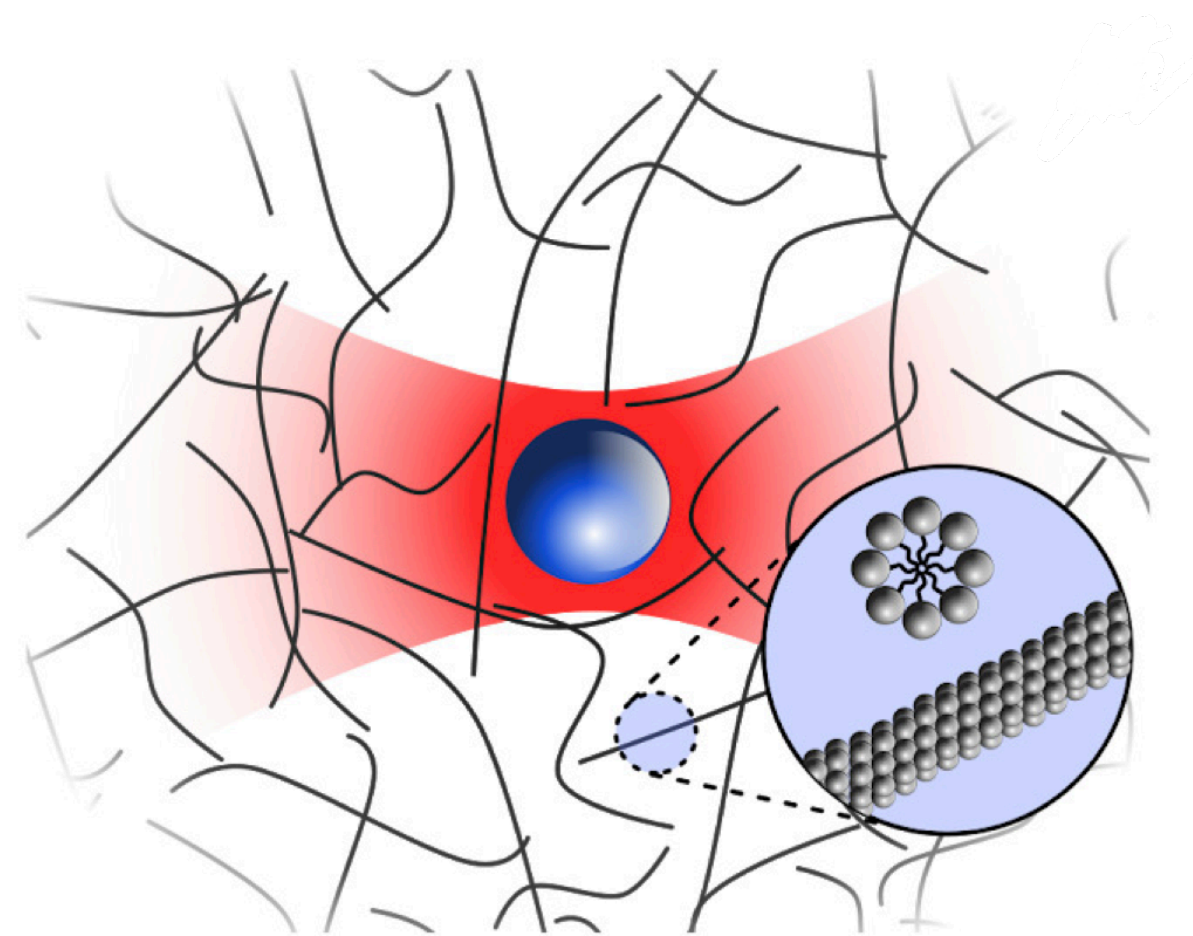
- **Brownian motion**

$$m \ddot{x}(t) = -\gamma \dot{x}(t) + \zeta(t) - V'(x(t))$$

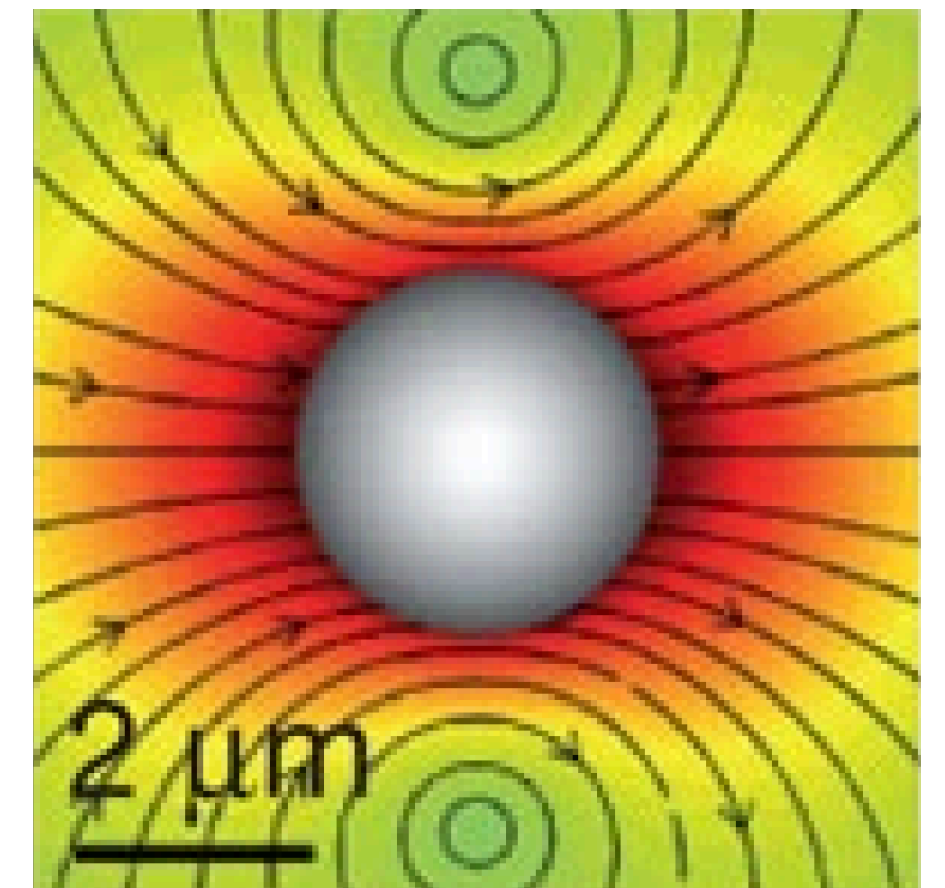
$$\langle \zeta(t) \zeta(t') \rangle = 2 k_B T \gamma \delta(t - t')$$

**Separation of timescales
challenged by, e.g.,**

- Viscoelastic fluids
- Near-critical media
- Hydrodynamic memory



[R. Jain et al., J. Chem. Phys. (2021)]



[T. Franosh et al., Nature (2011)]

Particle in a complex medium

From Brownian motion to correlated media

- **Brownian motion**

$$m \ddot{x}(t) = -\gamma \dot{x}(t) + \zeta(t) - V'(x(t))$$

$$\langle \zeta(t) \zeta(t') \rangle = 2 k_B T \gamma \delta(t - t')$$

- **GLE**

$$m \ddot{x}(t) = - \int^t dt' \Gamma(t - t') \dot{x}(t') + \zeta(t) - V'(x(t))$$

$$\langle \zeta(t) \zeta(t') \rangle = k_B T \Gamma(|t - t'|)$$

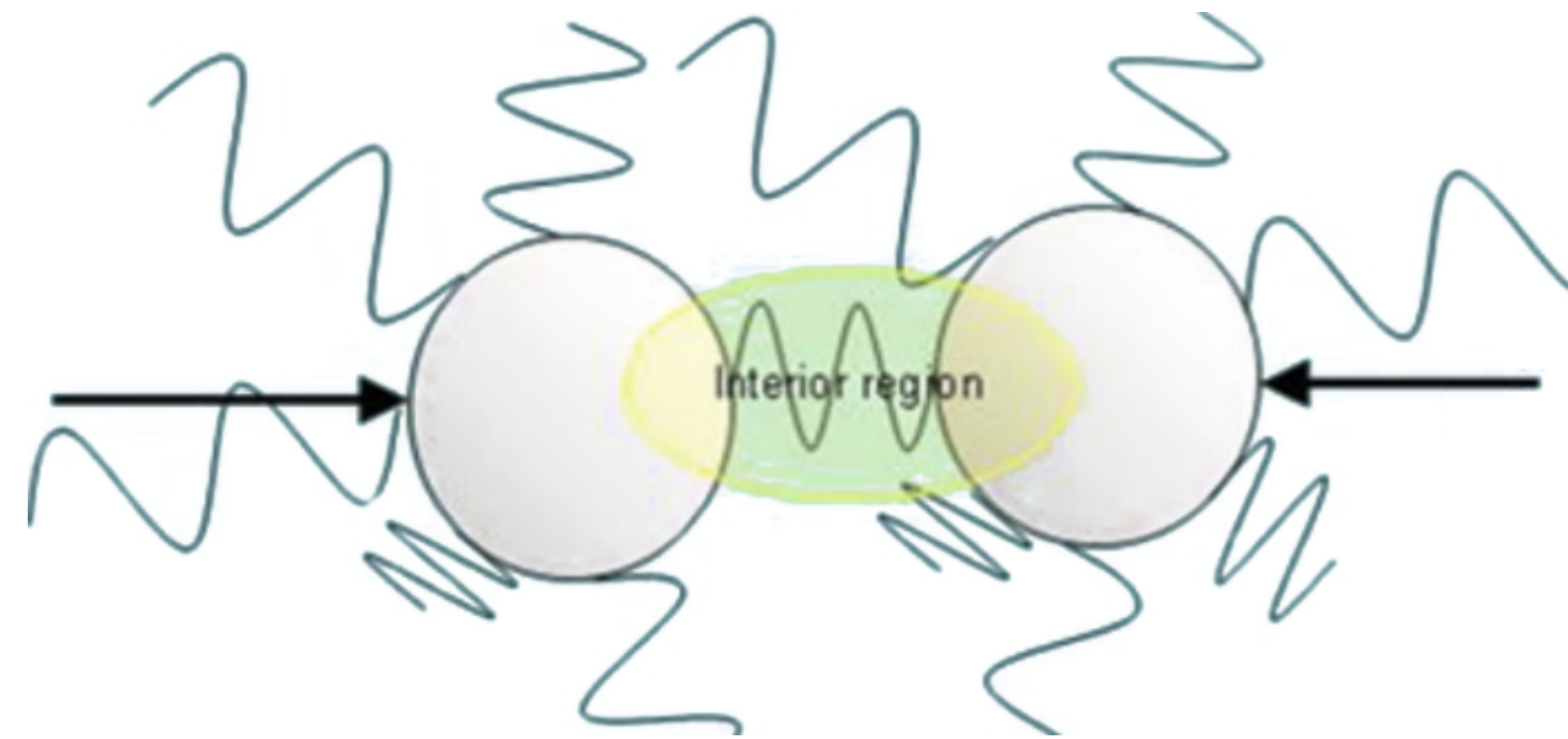
[e.g. Caldeira&Leggett '83]



Spatial structure? Energy flows? Fluctuation-induced forces?

Fluctuation-induced forces

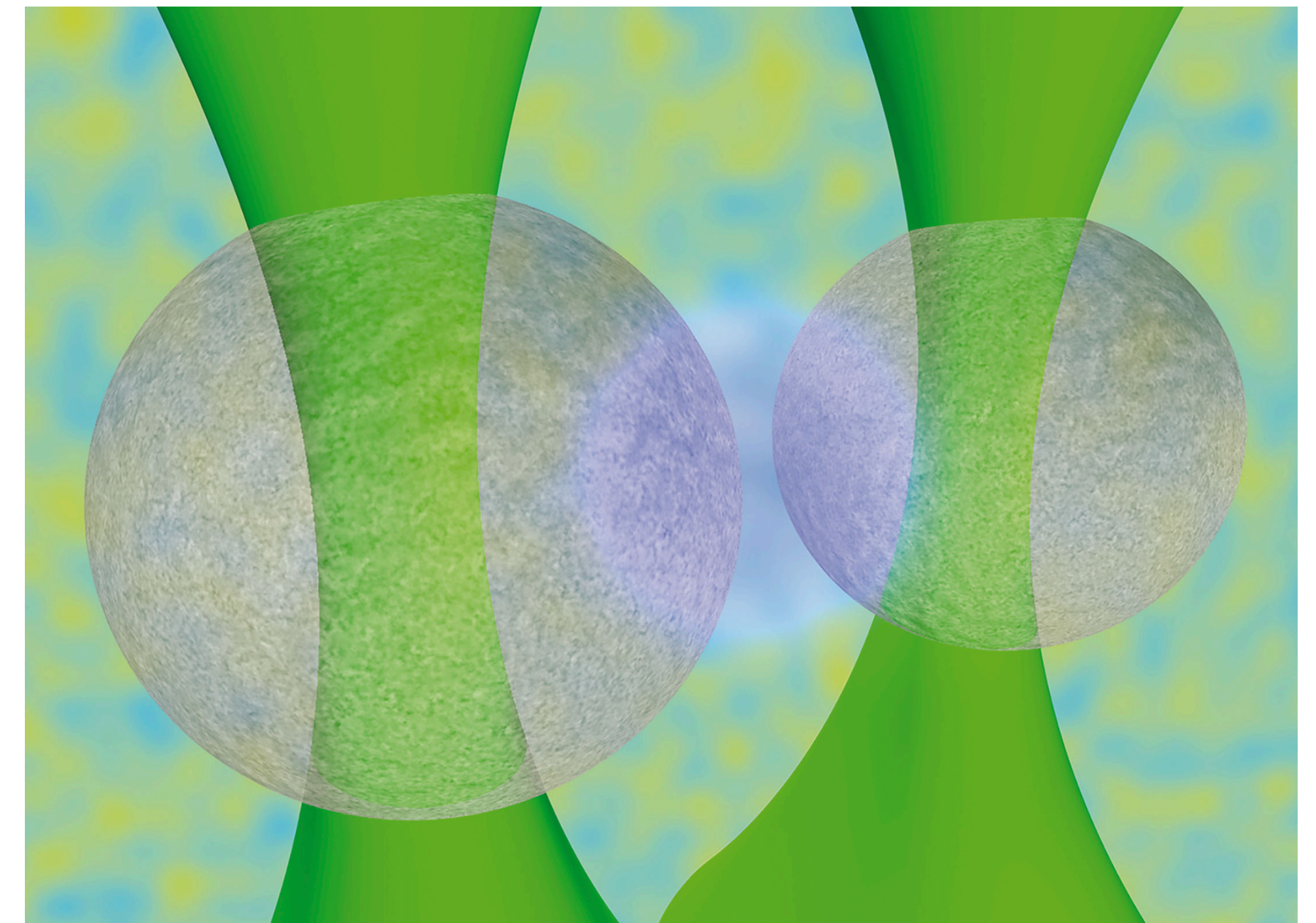
in correlated systems



- Critical Casimir forces
- Near-critical media (e.g., binary liquid mixtures)
- Retardation effects (not instantaneous)

[C. Hertlein et al., Nature (2008)]

[A. Magazzù et al., Soft Matter (2018)]



A toy model

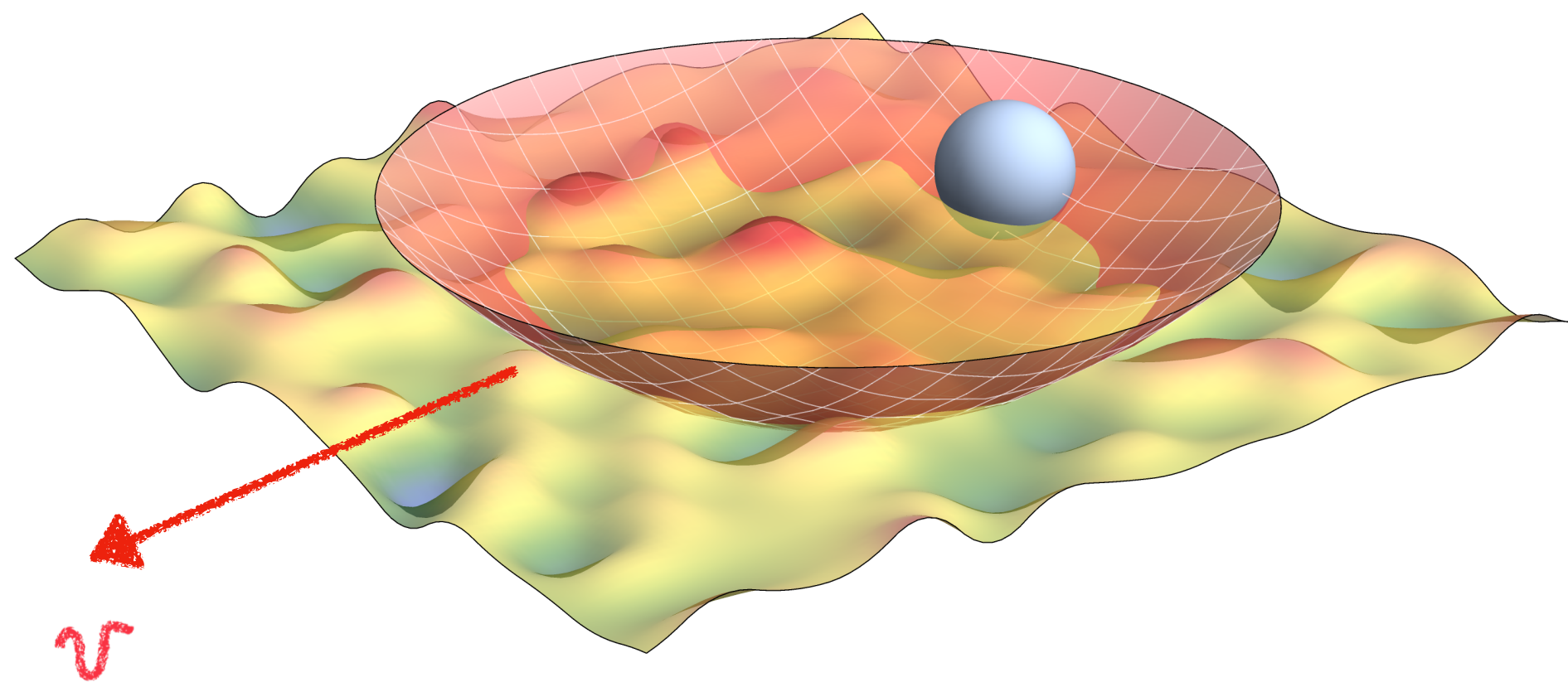
$\xi = r^{-1/2}$ sets the range of spatial correlations of $\phi(x, t)$

$$\mathcal{H}[\phi, \mathbf{Y}] = \mathcal{H}_\phi[\phi] + \mathcal{U}(\mathbf{Y}, t) + \mathcal{H}_{\text{int}}[\phi, \mathbf{Y}]$$

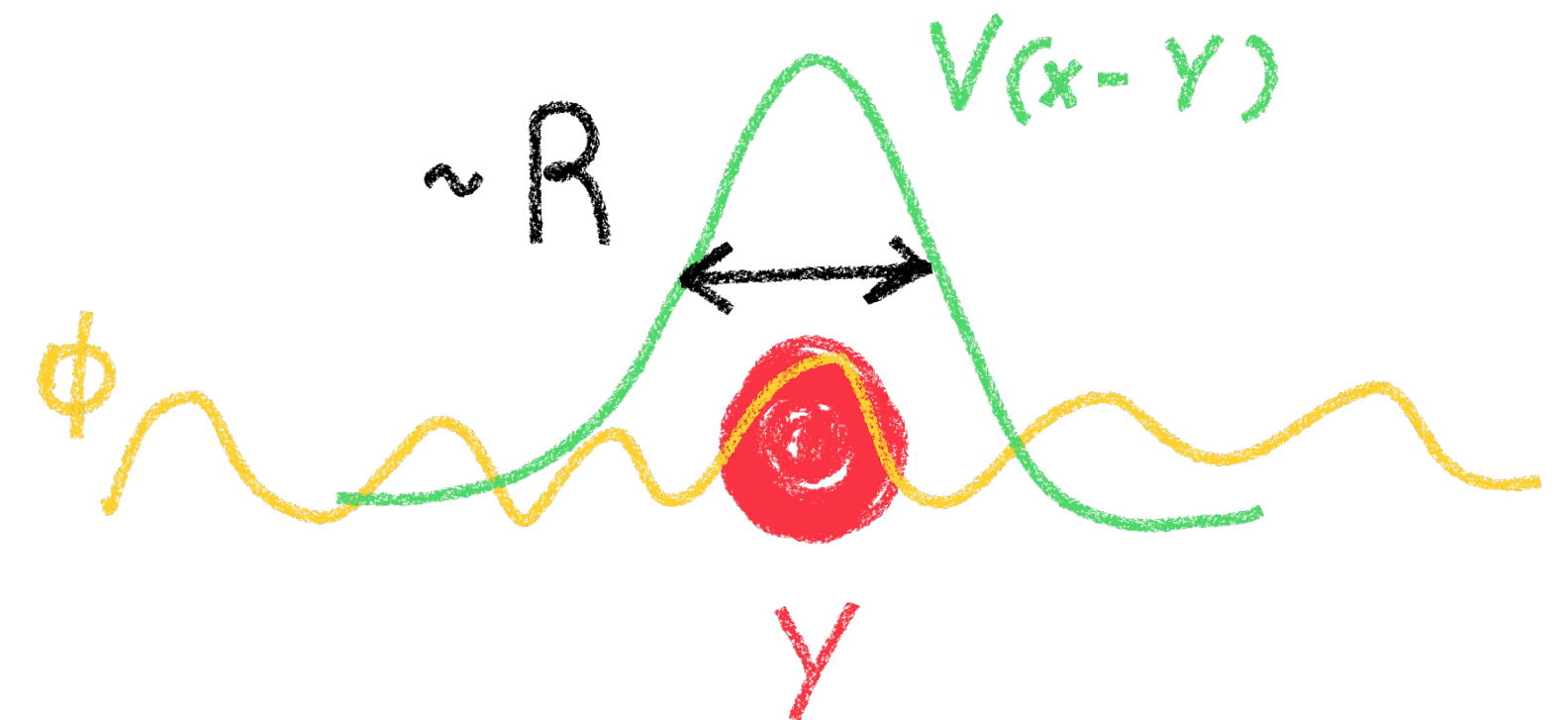
$$\mathcal{H}_\phi = \int d^d \mathbf{x} \left[\frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} r \phi^2 \right]$$

$$U(\mathbf{Y}, t) = \frac{\kappa}{2} (\mathbf{Y} - \mathbf{v}t)^2$$

$$\mathcal{H}_{\text{int}} = -\lambda \int d^d x V(\mathbf{x} - \mathbf{Y}) \phi(\mathbf{x})$$



$V(\mathbf{x} - \mathbf{Y})$ extends within the size R of the particle



Dynamics

ϕ , $\mathbf{Y}$ influence each other:

$$\gamma_y \dot{\mathbf{Y}}(t) = -\nabla_{\mathbf{Y}} \mathcal{H} + \boldsymbol{\xi}(t)$$

$$\gamma_\phi \dot{\phi}(\mathbf{x}, t) = -(-\nabla^2)^{\alpha/2} \frac{\delta \mathcal{H}}{\delta \phi(\mathbf{x}, t)} + \eta(\mathbf{x}, t)$$

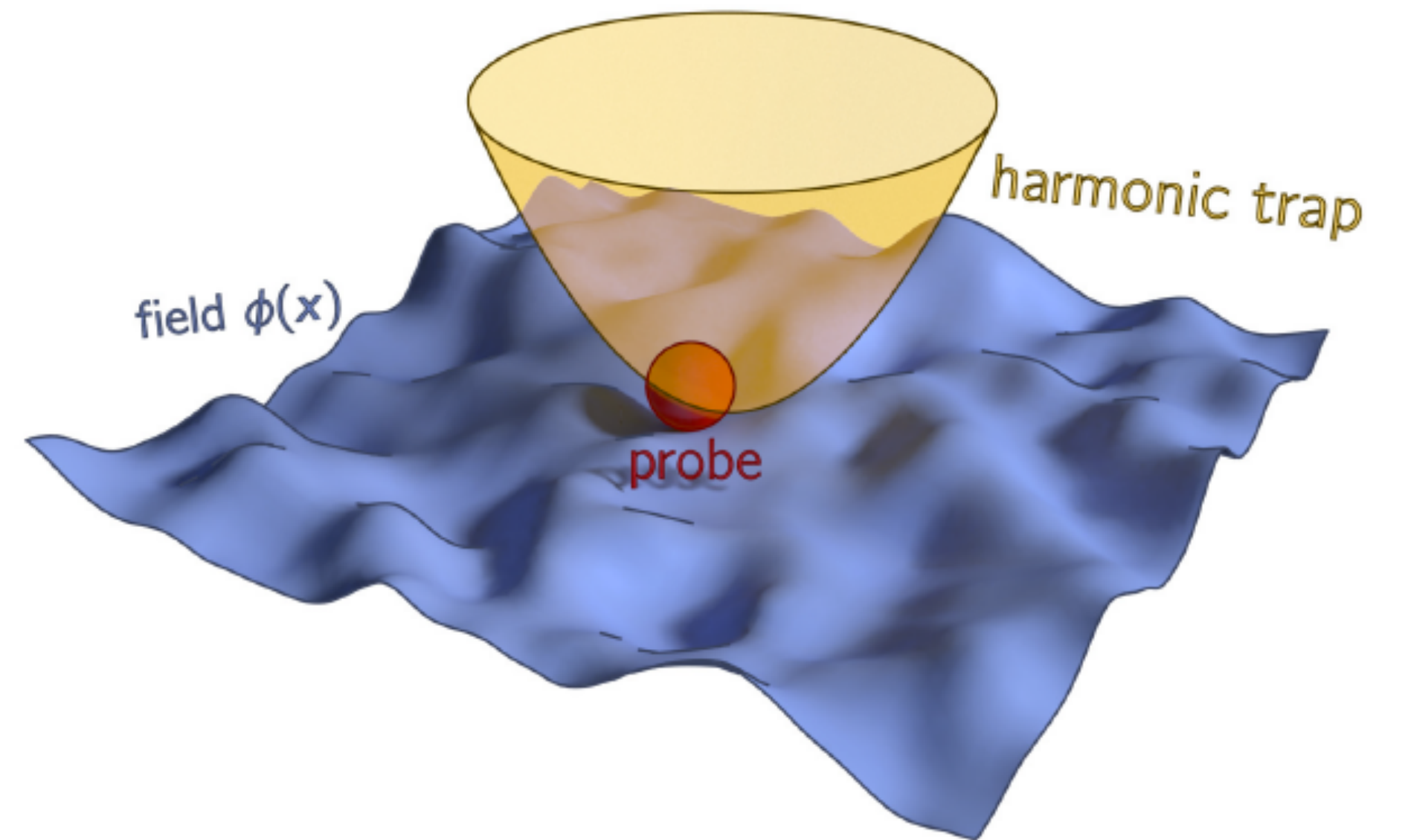
+ thermal bath @T,

$$\left\{ \begin{array}{l} \langle \xi_i(t) \xi_j(t') \rangle = 2\gamma_y T \delta_{ij} \delta(t - t') \\ \langle \eta(\mathbf{x}, t) \eta(\mathbf{x}', t') \rangle = 2\gamma_\phi T (-\nabla^2)^{\alpha/2} \delta^d(\mathbf{x} - \mathbf{x}') \delta(t - t') \end{array} \right.$$

Go in Fourier:

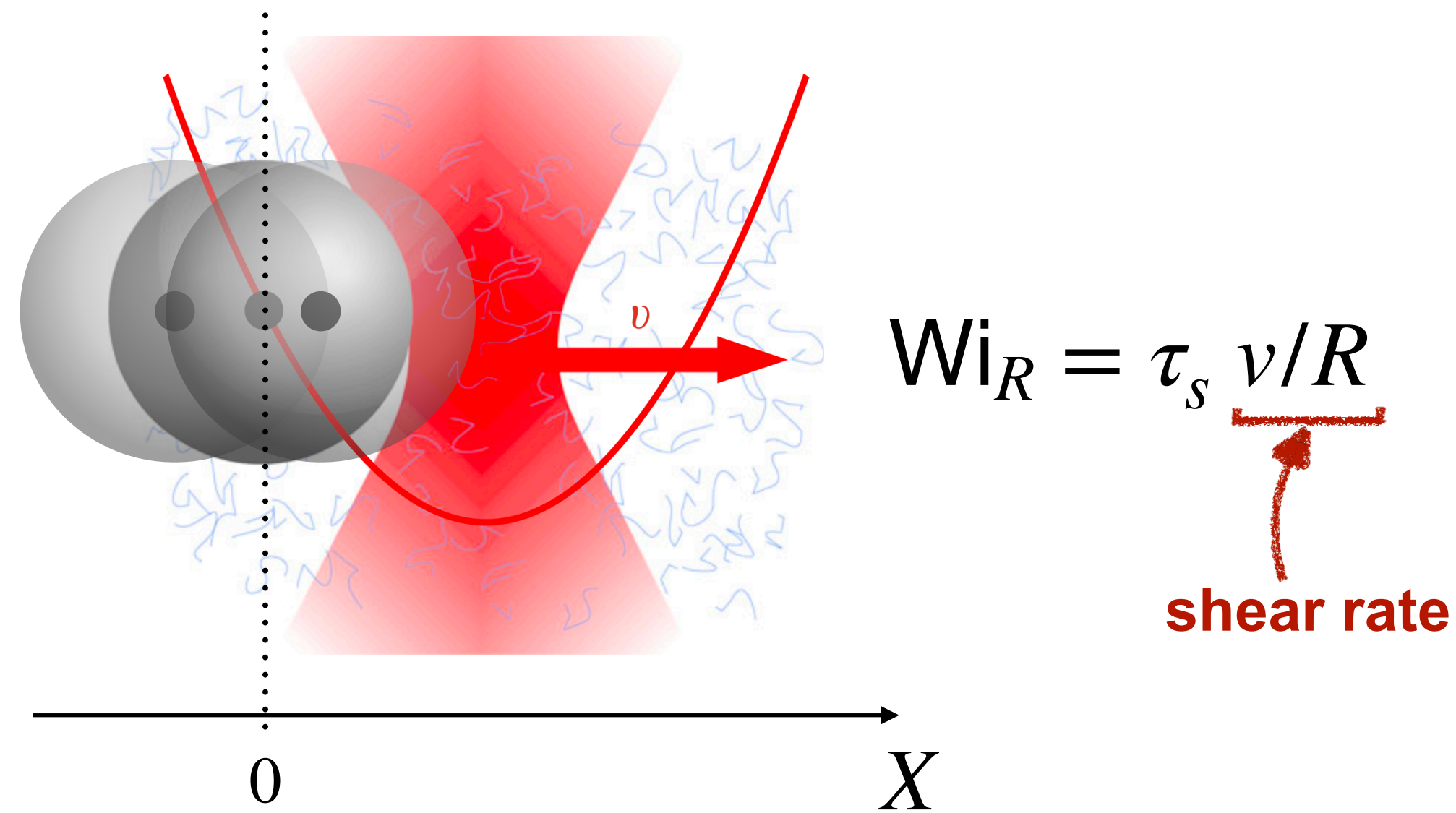
$$\gamma_y \dot{\mathbf{Y}} = -\kappa(\mathbf{Y} - \mathbf{v}t) + \lambda \int_{\mathbb{R}} \frac{d^d q}{(2\pi)^d} i\mathbf{q} V_{-q} \phi_q e^{i\mathbf{q} \cdot \mathbf{Y}} + \boldsymbol{\xi}$$

$$\gamma_\phi \dot{\phi}_q = -q^\alpha (q^2 + r) \phi_q + \cancel{\lambda} q^\alpha V_q e^{-i\mathbf{q} \cdot \mathbf{Y}} + \eta_q \quad \longrightarrow \quad \tau_\phi^{-1}(q) \equiv \gamma_\phi^{-1} q^\alpha (q^2 + r)$$

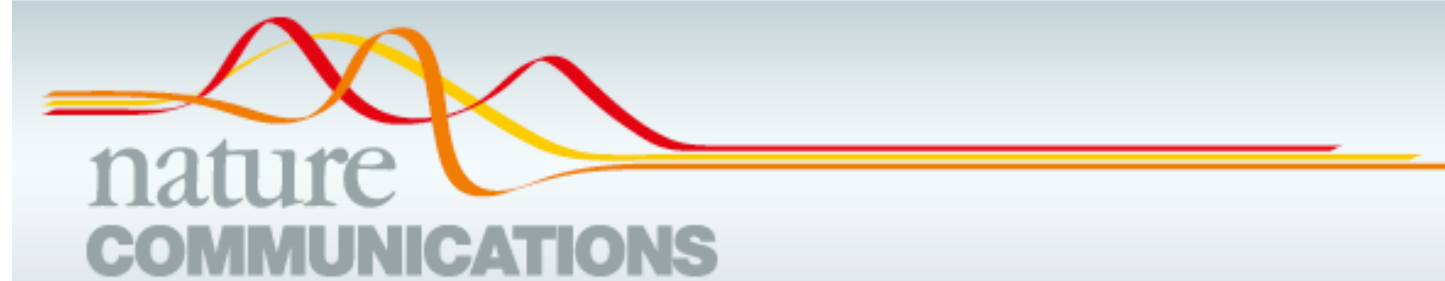


Memory-induced oscillations

in viscoelastic media



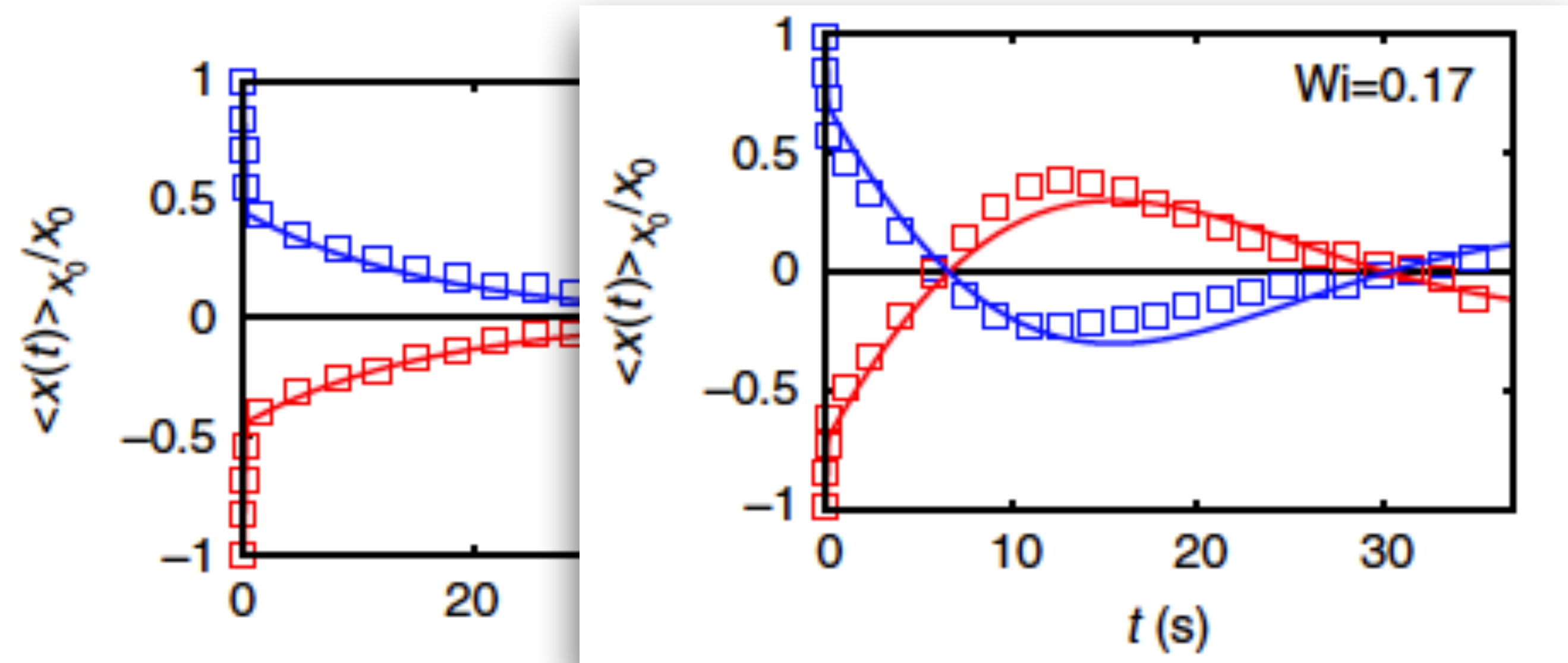
Micellar solution (overdamped!)



[2018]

Oscillating modes of driven colloids in overdamped systems

Johannes Berner^{1,2}, Boris Müller^{3,4}, Juan Ruben Gomez-Solano^{1,2}, Matthias Krüger^{3,4,5} & Clemens Bechinger^{1,2,4}



Memory-induced oscillations in correlated media

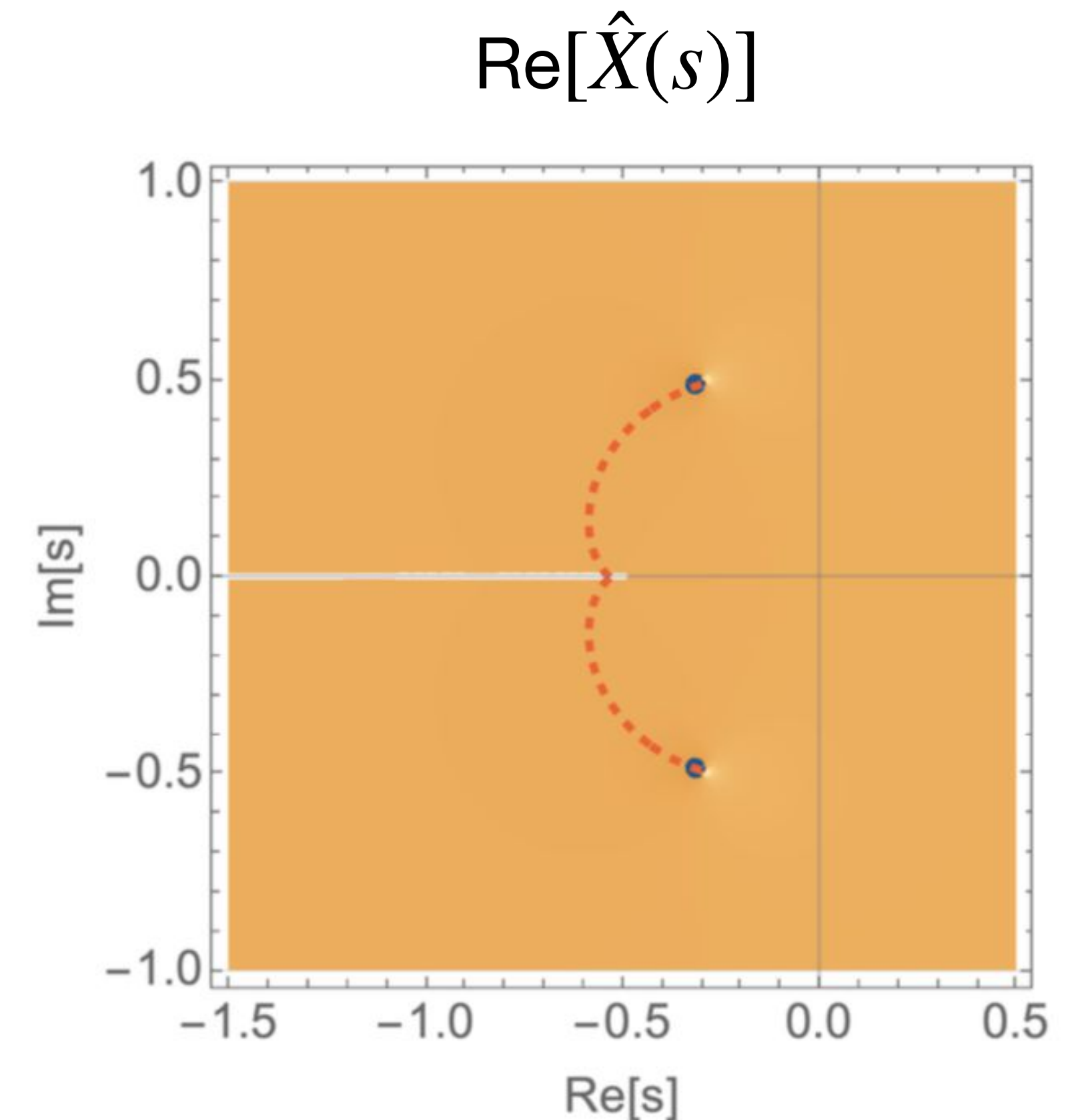
Linearized effective dynamics:

$$\partial_t \langle X_j(t) \rangle = -\langle X_j(t) \rangle [\gamma + \hat{\Gamma}_j(s=0)] + \int_{-\infty}^t du \Gamma_j(t-u) \langle X_j(u) \rangle$$

$$\langle \hat{X}_j(s) \rangle = \frac{X_0}{s + \gamma - [\hat{\Gamma}_j(s) - \hat{\Gamma}_j(0)]}$$

Complex poles for sufficiently large:

- Field-particle coupling
- $Wi_R = \tau_R \nu / R$, with $\tau_R \equiv \tau_\phi(q \sim \pi/R)$



Memory-induced oscillations in correlated media

Linearized effective dynamics:

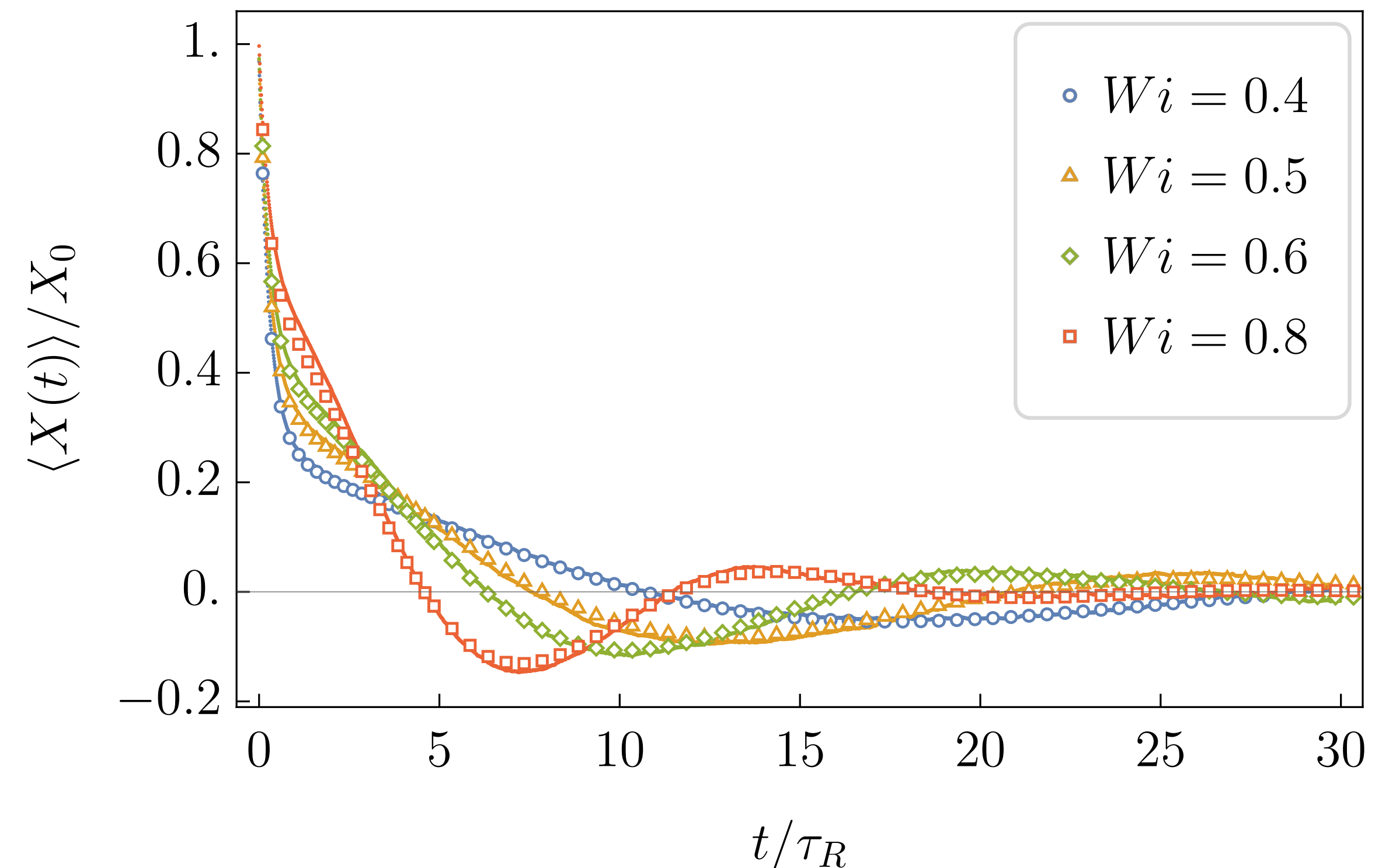
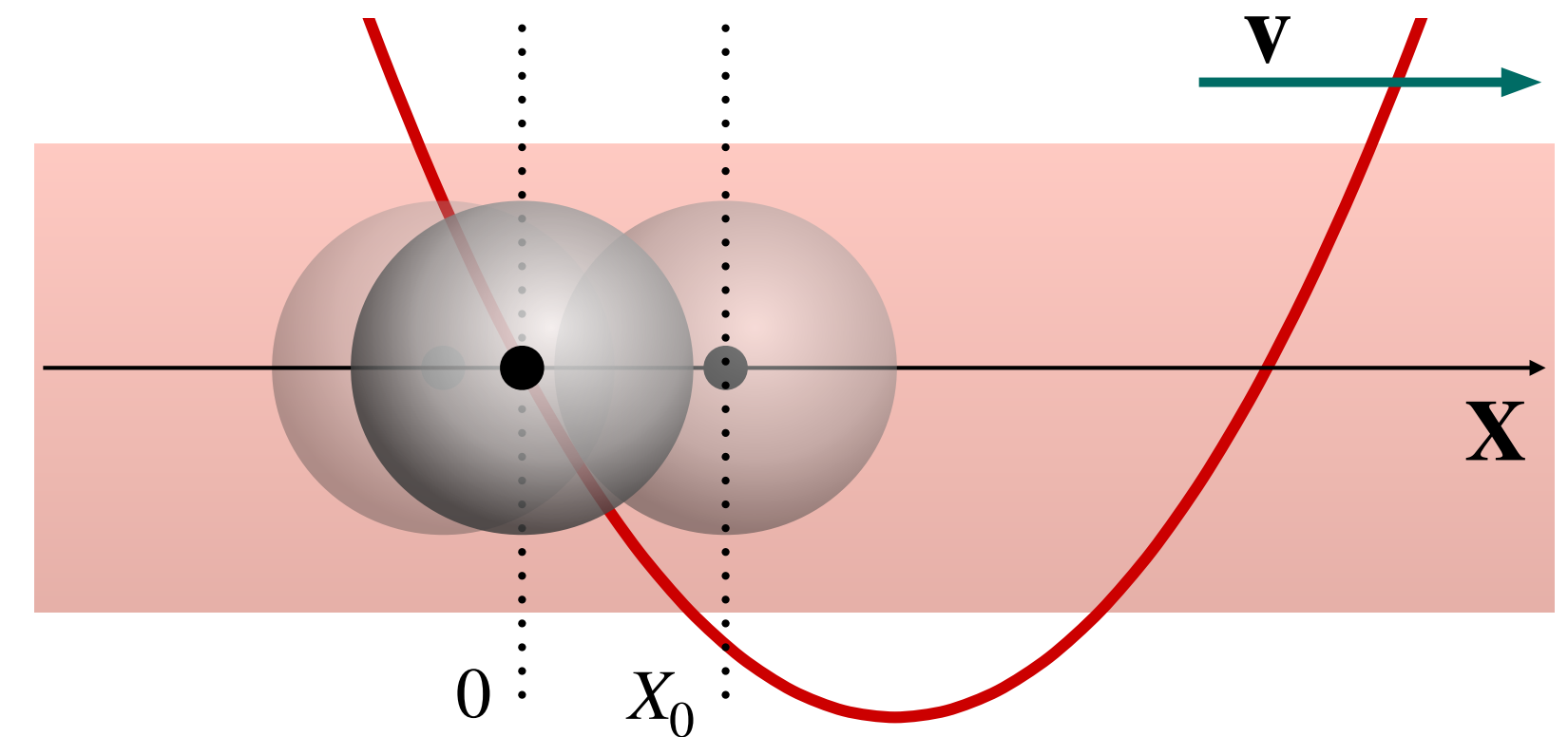
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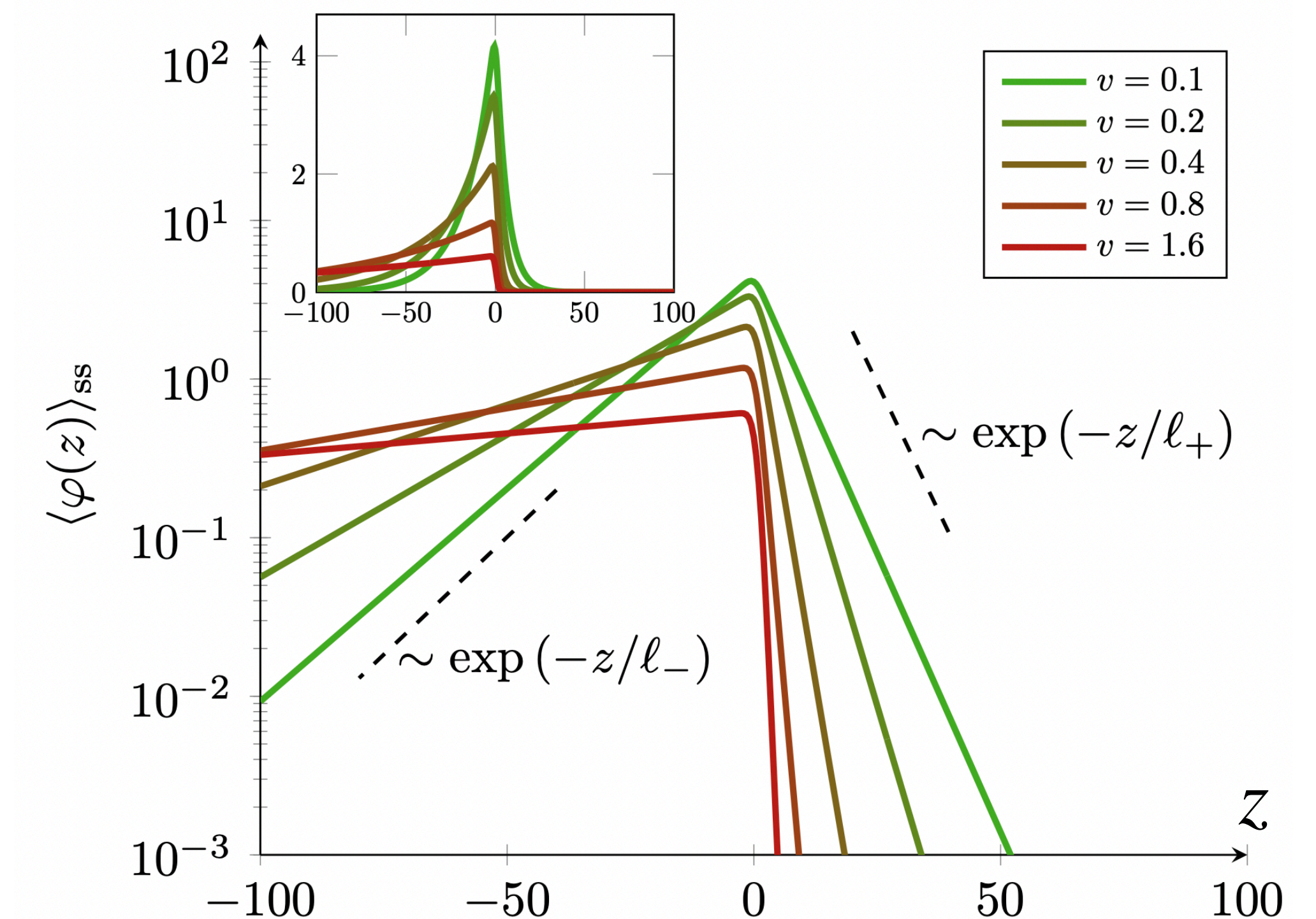
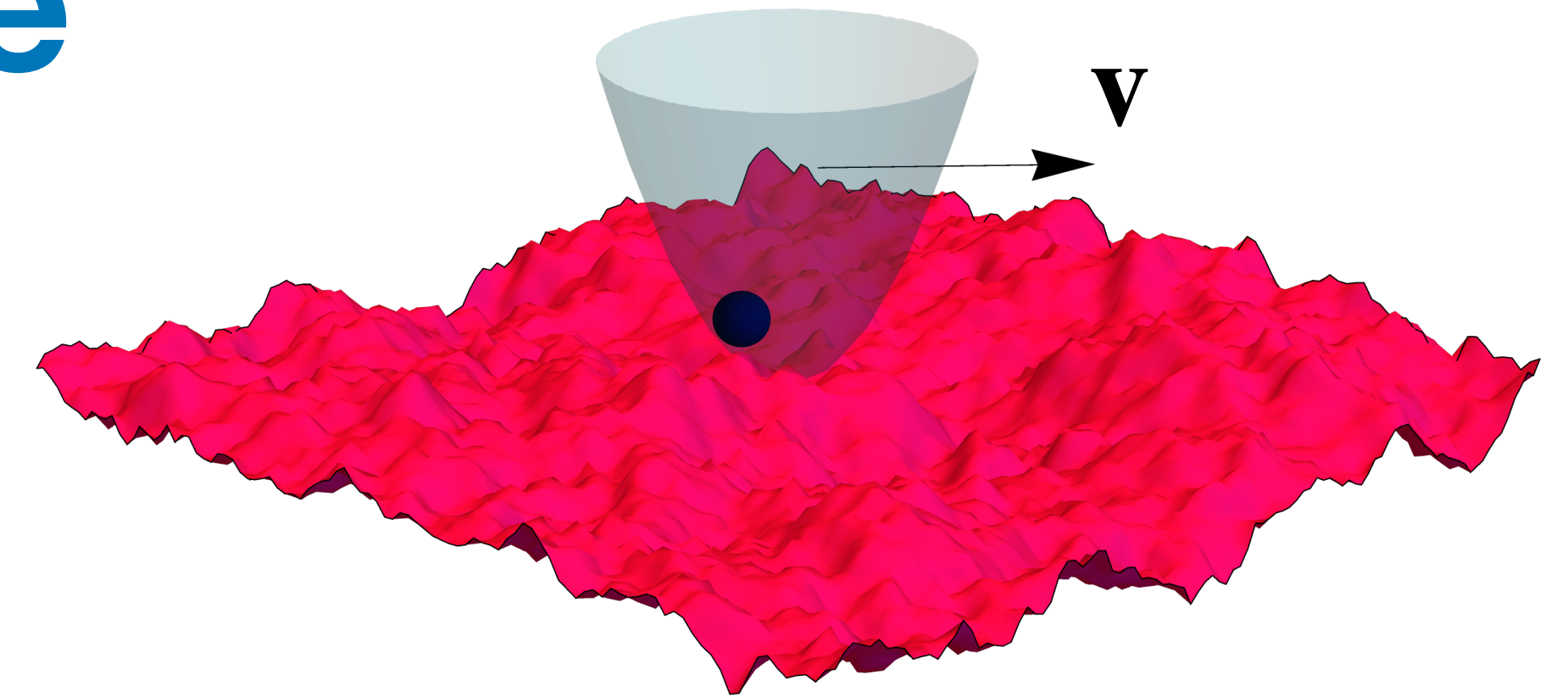
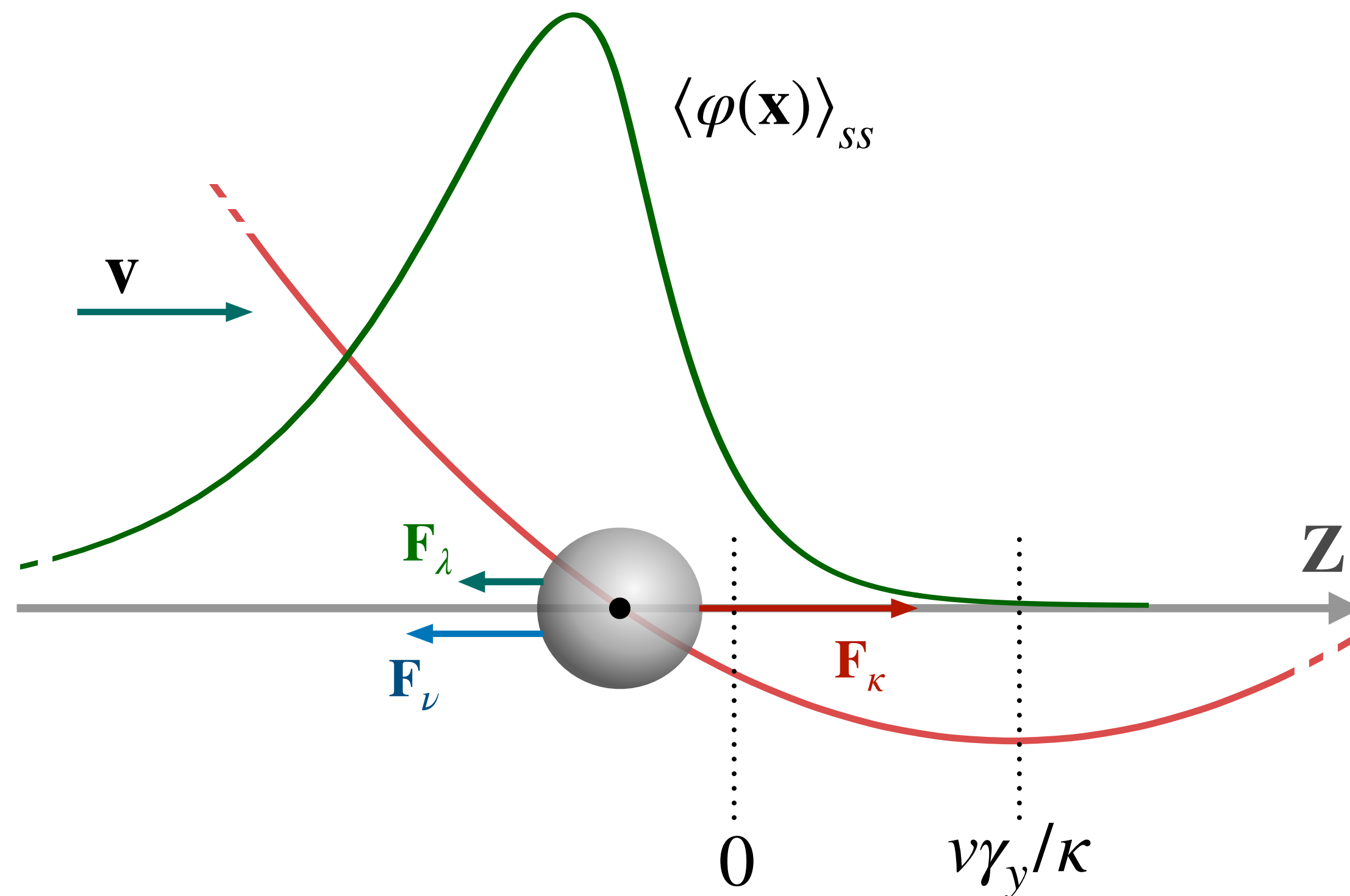
[DV, A.Gambassi, arXiv:2304.09684]



Nonequilibrium steady state

Particle driven at constant velocity

- $\mathbf{Z}$ = particle position in comoving frame
- $\langle \varphi \rangle_{ss}$ = avg field in comoving frame (*shadow*)



Stochastic thermodynamics

First law

$$d\mathcal{U} = dQ_y + dW_y^{\text{int}} + dW$$

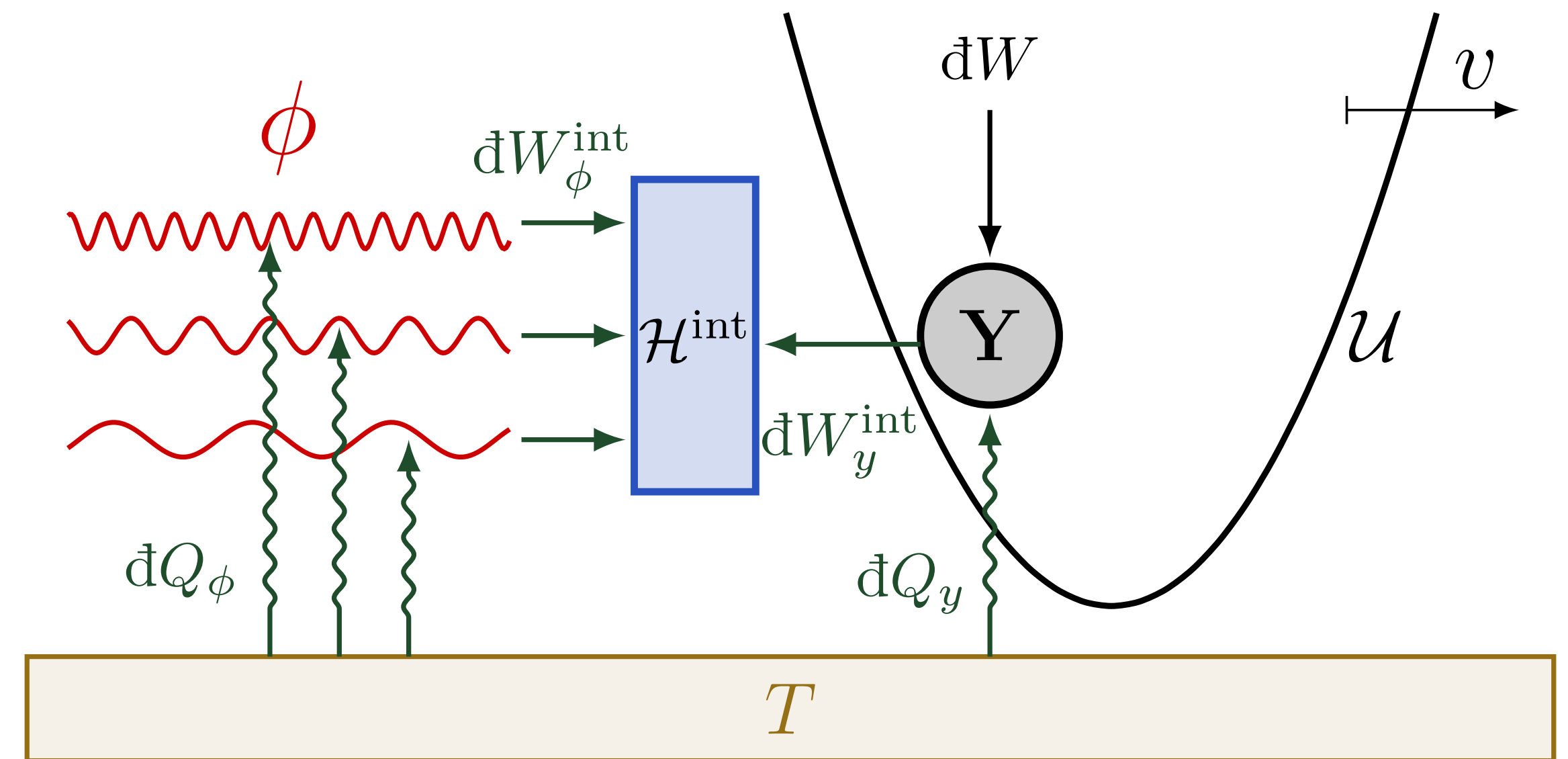
$$d\mathcal{H}_\phi = \int d^d \mathbf{x} \left[dQ_\phi(\mathbf{x}) + dW_\phi^{\text{int}}(\mathbf{x}) \right]$$



$$dW_y^{\text{int}} := -\nabla_Y \mathcal{H}^{\text{int}} \circ dY,$$

$$dW_\phi^{\text{int}}(\mathbf{x}) := -\frac{\delta \mathcal{H}^{\text{int}}}{\delta \phi(\mathbf{x})} \circ d\phi(\mathbf{x}).$$

$$\mathcal{H}[\phi, Y, t] = \mathcal{H}_\phi[\phi] + \mathcal{H}^{\text{int}}[\phi, Y] + \mathcal{U}(Y, t)$$



[DV, S. M. Loos, B. Walter, É. Roldán, A. Gambassi, arxiv:2305.16235]

Stochastic thermodynamics

Average dissipated power

- Full CGF of the dissipated **power**

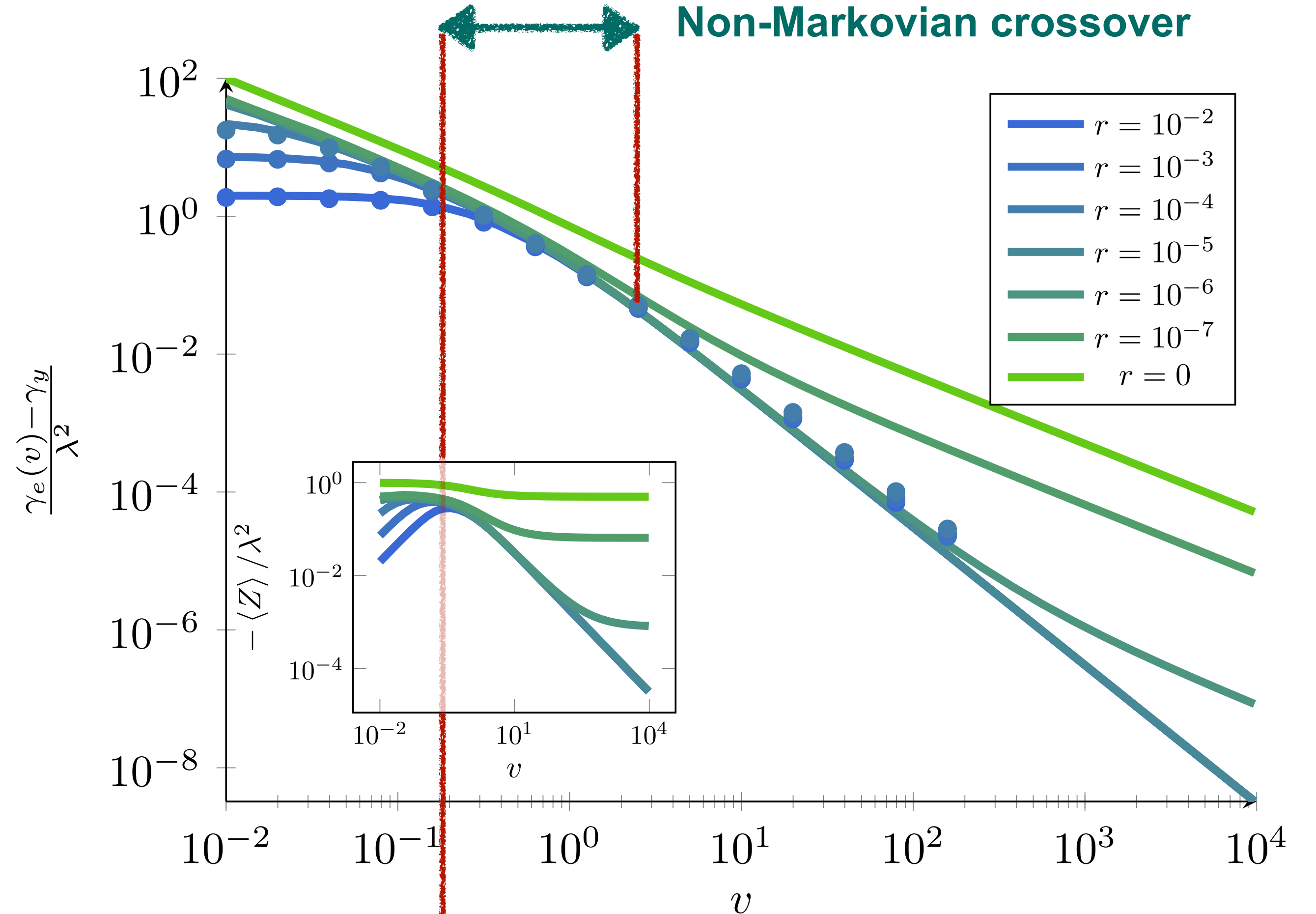
$$\dot{W} = -\kappa \mathbf{v} \cdot (\mathbf{Y} - \mathbf{v}t)$$

- Average value

$$\langle \dot{W} \rangle \equiv \gamma_e(v) v^2$$

- In the uncoupled case,

$$\gamma_e(v) \equiv \gamma_y$$



$$\text{Wi}_\xi = \tau_\phi(q \sim \pi/\xi) v/\xi \equiv 1$$

Stochastic thermodynamics

Second law

$$\Delta S_{\text{tot}}[\{\mathbf{Y}, \phi\}_{t_i}^{t_f}] = \ln \frac{\mathcal{P}[\{\mathbf{Y}, \phi\}_{t_i}^{t_f}]}{\mathcal{P}^R[\{\mathbf{Y}^R, \phi^R\}_{t_i}^{t_f}]} \geq 0$$

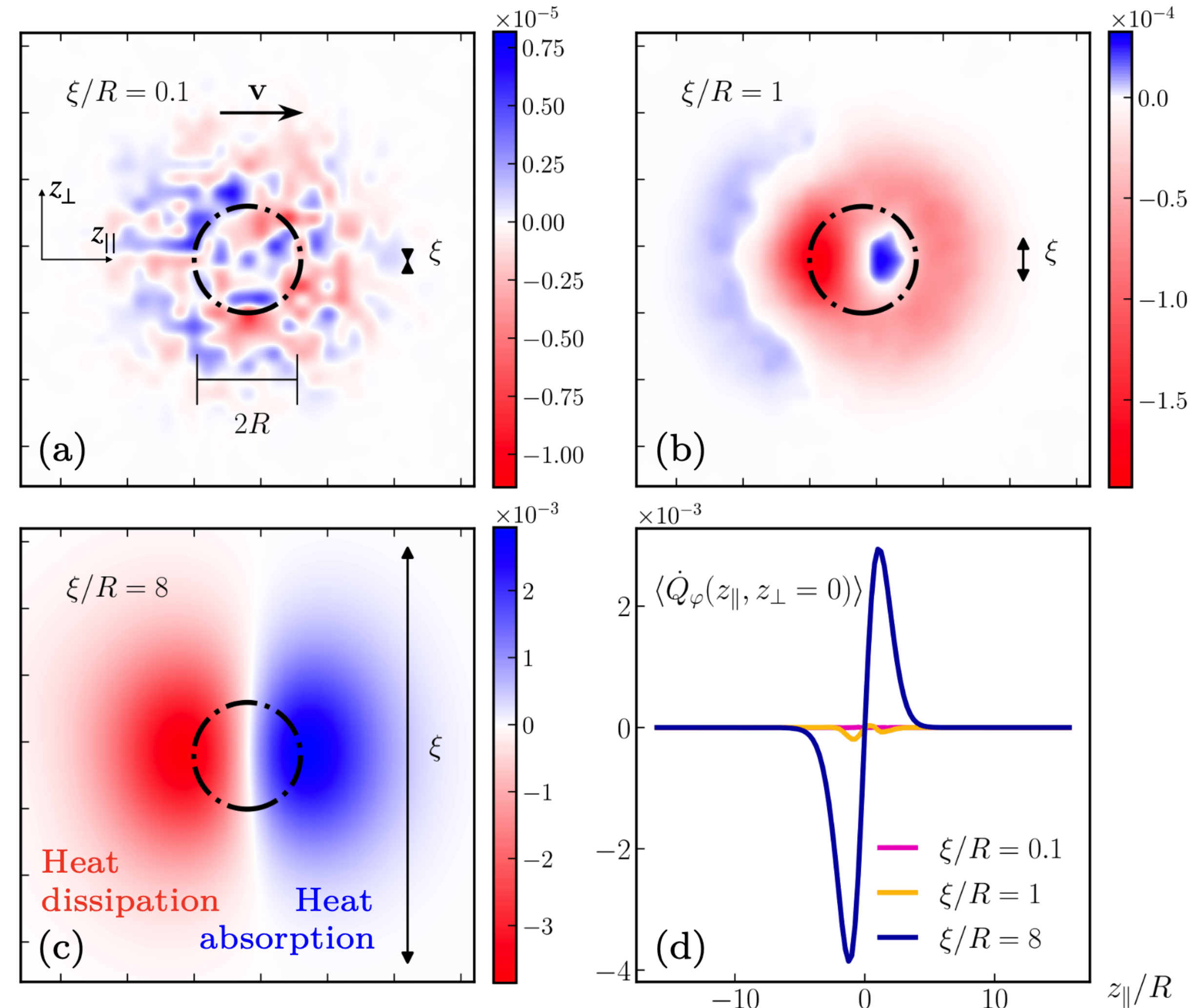
$$= -\frac{Q_y}{T} - \int d^d \mathbf{x} \frac{Q_\phi(\mathbf{x})}{T} + \Delta S_{y,\phi}^{\text{sh}}$$

with

$$S_{y,\phi}^{\text{sh}} \propto -\ln \rho_{\mathbf{Y},\phi} \quad \leftarrow \text{Shannon entropy}$$

$$\mathring{d}Q_y = (\xi - \gamma_y \dot{\mathbf{Y}}) \circ d\mathbf{Y}$$

$$\mathring{d}Q_\phi(\mathbf{x}) = \left\{ (-\nabla^{-2})^{\alpha/2} [\eta(\mathbf{x}) - \gamma_\phi \dot{\phi}(\mathbf{x})] \right\} \circ d\phi(\mathbf{x})$$



To sum up

- Particle + field with spatio-temporal correlations
- Memory-induced oscillations
- Spatially-resolved stochastic thermodynamics

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